

PH241 Final Solutions

Prob 1 (10pt).

(a) (5pt)

1) $3d$ to $2p$ transition, energy difference is

$$E = (-13.6 \text{ eV}) \left(\frac{1}{3^2} - \frac{1}{2^2} \right) = 1.889 \text{ eV}$$

$$\lambda = \frac{hc}{E} = 656.11 \text{ nm} \dots\dots +2\text{pts}$$

$$\Delta\lambda = \frac{\lambda^2}{hc} \Delta E = 0.08 \text{ nm} \dots\dots +1\text{pt}$$

So, wavelength of the three normal Zeeman components,

656.03 nm , 656.11 nm , 656.19 nm $\dots\dots +1\text{pt}$

2) Since E only depends on n , $3s$ to $2p$ transition is same as 1) $\dots\dots +1\text{pt}$

(b) (5pt) There are 6 possible sets for each $2p$ electrons. $\dots\dots +2\text{pts}$ So, $6 \times 6 = 36$ possibilities. $\dots\dots +1\text{pt}$ Also, there are 6 combinations in which the two sets are identical. So, by Pauli principle, $36 - 6 = 30$ possibilities. $\dots\dots +2\text{pts}$

Prob 2 (10pt).

(a) (5pt)

$$k = \frac{2(E - E_{\min})}{(R - R_{eq})^2} \simeq 310 \text{ eV/nm}^2$$

Also, from $\mu = \frac{m(\text{Na})m(\text{Cl})}{m(\text{Na}) + m(\text{Cl})} \simeq 13.96u$, $\dots\dots +2\text{pts}$

$$f = \frac{1}{2\pi} \sqrt{\frac{k}{\mu}} \simeq 7.4 \times 10^{12} \text{ Hz}$$

So, $\lambda = \frac{c}{f} \simeq 40 \mu\text{m}$ (infrared region) $\dots\dots +1\text{pt}$ and $E = hf \simeq 0.031 \text{ eV}$. $\dots\dots +1\text{pt}$

Since it's valid about 0.25 eV , maximum vibrational quantum number is 8. $\dots\dots +1\text{pt}$

(b) From $\mu = \frac{m(\text{Na})m(\text{Cl})}{m(\text{Na}) + m(\text{Cl})} \simeq 13.96u$ and $R_{eq} = 0.236 \text{ nm}$, we can get

$$B = \frac{\hbar^2}{2mR_{eq}^2} \simeq 2.68 \times 10^{-5} \text{ eV} \dots\dots +2\text{pts}$$

So,

$$L=1 \text{ to } L=0 \text{ (lowest)} : \Delta E = 2B = 5.36 \times 10^{-5} \text{ eV} \quad \& \quad \lambda = \frac{hc}{\Delta E} \simeq 23.1 \text{ nm} \dots\dots$$

$+2\text{pts}$

$$f = \frac{c}{\lambda} \simeq 1.3 \times 10^{10} \text{ Hz (Microwave region)} \dots\dots +1\text{pt}$$

Prob 3 (10pt).

(a) (5pt) $T = 1000\text{K} \Rightarrow kT \simeq 0.0862 \text{ eV} \dots\dots +1\text{pt}$

So, $\frac{N_2}{N_1} = \frac{d_2 e^{-E_2/kT}}{d_1 e^{-E_1/kT}} = 3 e^{-(E_2-E_1)/kT} \simeq 0.295 \dots\dots +2\text{pts}$

With, $N = N_1 + N_2$, we can get $N_1 \simeq 0.77N$, $N_2 \simeq 0.23N$. $\dots\dots +2\text{pts}$

(b) (5pt) $T = 300\text{K} \Rightarrow kT \simeq 0.02585 \text{ eV} \dots\dots +1\text{pt}$

$E - E_F = \frac{E_g}{2}$ (by problem assumption) $\dots\dots +1\text{pt}$

So,

$$f_{FD}(E) = \frac{1}{e^{(E-E_F)/kT} + 1} = \frac{1}{e^{E_g/2kT} + 1} \simeq 6.4 \times 10^{-47} \text{ for C}$$

$$f_{FD}(E) = \frac{1}{e^{(E-E_F)/kT} + 1} = \frac{1}{e^{E_g/2kT} + 1} \simeq 5.8 \times 10^{-10} \text{ for Si.} \dots\dots +2\text{pts}$$

The ratio of electron concentration is 1.1×10^{-37} . $\dots\dots +1\text{pt}$

Prob 4 (10pt).

(a) (5pt) Maxwell distribution : distribution of velocity (i.e., $f(v) = C_1 e^{-\frac{\frac{1}{2}mv^2}{kT}} \dots\dots +2\text{pts}$)

Boltzmann distribution : distribution of energy (i.e., $f(E) = C_2 e^{-\frac{E}{kT}} \dots\dots +2\text{pts}$)
 $+1\text{pt for mention the velocity \& energy}$

(b) (5pt) For number density distribution of energy $n(E)$,

for collision $1 + 2 \rightarrow 3 + 4$, rate of collision $R(1 + 2 \rightarrow 3 + 4) = C_1 n(E_1)n(E_2)$ is satisfied.

And for collision $3 + 4 \rightarrow 1 + 2$, $R(3 + 4 \rightarrow 1 + 2) = C_2 n(E_3)n(E_4)$ is satisfied. $\dots\dots +1\text{pt}$

From micro-reversibility, we can get $C_1 = C_2$, and so $n(E_1)n(E_2) = n(E_3)n(E_4)$. $\dots\dots +1\text{pt}$

Take the logarithm, we can get $\ln n(E_1) + \ln n(E_2) = \ln n(E_3) + \ln n(E_4)$.

With energy conservation $E_1 + E_2 = E_3 + E_4 \dots\dots +1\text{pt}$, we can guess $n(E)$ should be the form of $\ln n(E) = \alpha - \beta E$.

Equivalent to $n(E) = Ce^{-\beta E}$ +1pt

Since when $E = \frac{1}{2}mv^2$ (free particle energy), $n(E)$ should be same as Maxwell

distribution, so we can find that $\beta = \frac{1}{kT}$ +1pt

So, $n(E) = Ce^{-\frac{E}{kT}}$.

Prob 5 (10pt).

(a) (4pt) Mass of star $M = N \times m_n$ (m_n : mass of neutron)

$$E_{grav} = -\frac{3}{5} \frac{GM^2}{R} = -\frac{3}{5} \frac{GN^2 m_n^2}{R} \dots\dots +1pt$$

$$E_{neut} = NE_m = \frac{3}{5} NE_F = \frac{3}{5} N \frac{h^2}{2m_n} \left(\frac{3N}{8\pi V}\right)^{2/3} \dots\dots +1pt$$

$$E = E_{grav} + E_{neut}$$

Minimum radius occurs at $dE/dR = 0$, so, we can get the radius solving this :

$$R = \frac{h^2}{GN^{1/3} m_n^3} \left(\frac{9}{32\pi^2}\right)^{2/3} \dots\dots +2pts$$

(b) (4pt) $M_{sun} \approx 2.00 \times 10^{30} kg$, so $M = 4M_{sun} \approx 8.00 \times 10^{30} kg$

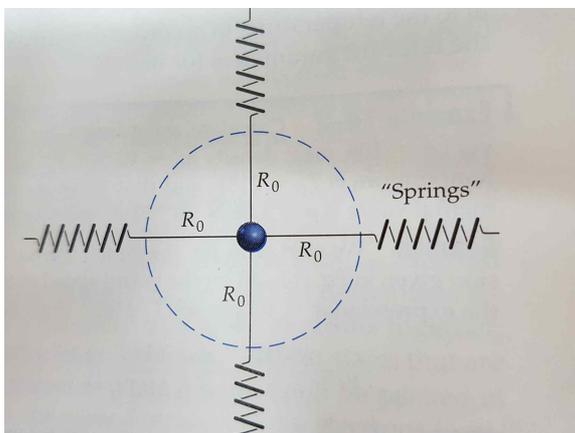
$$\& N = \frac{M}{m_n} \approx 4.8 \times 10^{57} \dots\dots +2pts$$

So, $R \approx 7.8 km$ +2pts

(c) (2pt) Density $\rho = \frac{M}{V} \approx 3.98 \times 10^{18} kg/m^3$ +2pts

Prob 6 (10pt).

(a) (5pt)



Let the amplitude R_0 , then the scattering cross section $\sigma = \pi R_0^2$ can be assumed. +1pt

Also, for this model, we can get kinetic energy $E_{kinetic} = \frac{1}{2}mv^2 = \frac{1}{2}mR_0^2\omega^2$ +1pt

Using the equipartition theorem, we get $\frac{1}{2}mR_0^2\omega^2 = \frac{1}{2}kT$ +1pt

So, $R_0^2 = \frac{kT}{mw^2}$.

And finally we get $\sigma = \pi \frac{kT}{mw^2}$ +2pts

(b) (5pt) Since $\sigma \propto T$ +1pt and other terms in ρ are temperature independent +2pts. So, we can easily check that $\rho \propto \sigma \propto T$ +2pts.

Prob 7 (10pt).

Heat capacity $C = C_{el} + C_{ph}$

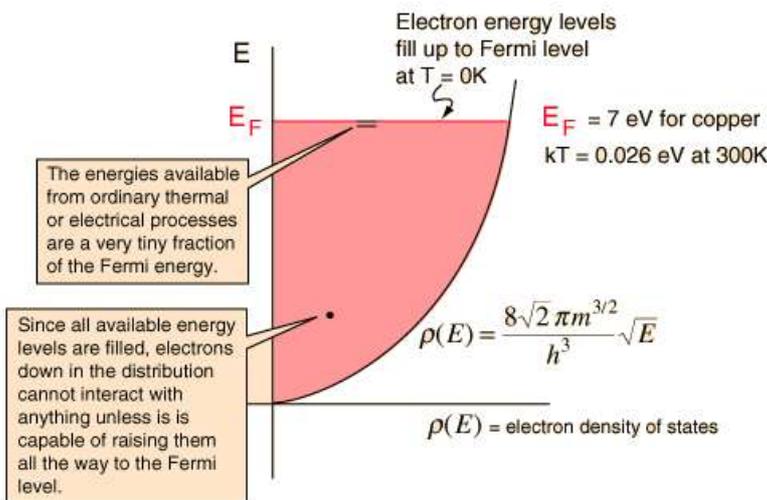
where C_{el} : contribution from electron and C_{ph} : contribution from phonon.

1) C_{el}

Using the Fermi-Dirac distribution ($N(E) = \frac{2}{e^{(E-E_F)/kT} + 1}$), we can find the

$\langle E \rangle = \frac{1}{N} \int EN(E)dE \propto T^2$.

Also, from the definition of heat capacity $C \equiv \frac{\partial U}{\partial T}$, we can get $C_{el} \propto \frac{T}{T_F}$.



For this C_{el} , since $k_B T$ (order of meV) is very small compared to E_F (order of few eV) of metals, only small number of electrons near E_F can contribute to C_{el} . So, it is very small.

2) C_{ph}

Since phonon is boson, quantum mechanically at low temperature, we have to use

Bose-Einstein statistics($f_{BE}(E) = \frac{1}{A_{BE}e^{E/kT} - 1}$). +3pts

Also, phonon has vibrational energy of $\hbar\omega_{ph}$ (order of 10meV) and at low temperature about 15K or below, $k_B T \sim 1meV$ or lower, so it should be

$$f_{BE}(E) \sim e^{-\frac{\hbar\omega_{ph}}{kT}}. \text{ +5pts}$$

[Using the Debye model(considering solid as a ensemble of independent of quantum harmonic oscillators), we can get $C_{ph} \propto (\frac{T}{T_D})^3$.]

So, at low temperature($\hbar\omega_{ph} \gg k_B T$), thermal energy is too small to excite the phonon and equipartition theorem can not be applied, phonon contribution becomes also very small.

So, for $C = C_{el} + C_{ph}$, C becomes very small when $T \rightarrow 0K$ +2pts.